

Optical Detectors

- **PURPOSE: TO CONVERT THE RECEIVED OPTICAL SIGNAL INTO AN ELECTRICAL SIGNAL.**

Requirements For Detector

- **HIGH SENSITIVITY** (at operating wave lengths) at normal op. temp (300 K) 0.85 μm /1.1 μm / 1.3 μm
- **HIGH FIDELITY** (Linear response over a wide range for analog transmission)
- **HIGH QUANTUM EFFICIENCY**
- **SHORT RESPONSE TIME.** (FOR SUITABLE BW)
- **MINIMUM NOISE** (Introduced by detector)

Optical Detectors (Contd.)

- **Stability of performance characteristics** (independent of change in ambient conditions).
- **SMALL SIZE.** (for efficient coupling to fiber and easy packaging)
- **LOW BIAS VOLTAGES.**
- **HIGH RELIABILITY**
- **LOW COST**

OPTICAL DETECTION - PRINCIPLE

Optical detectors 423

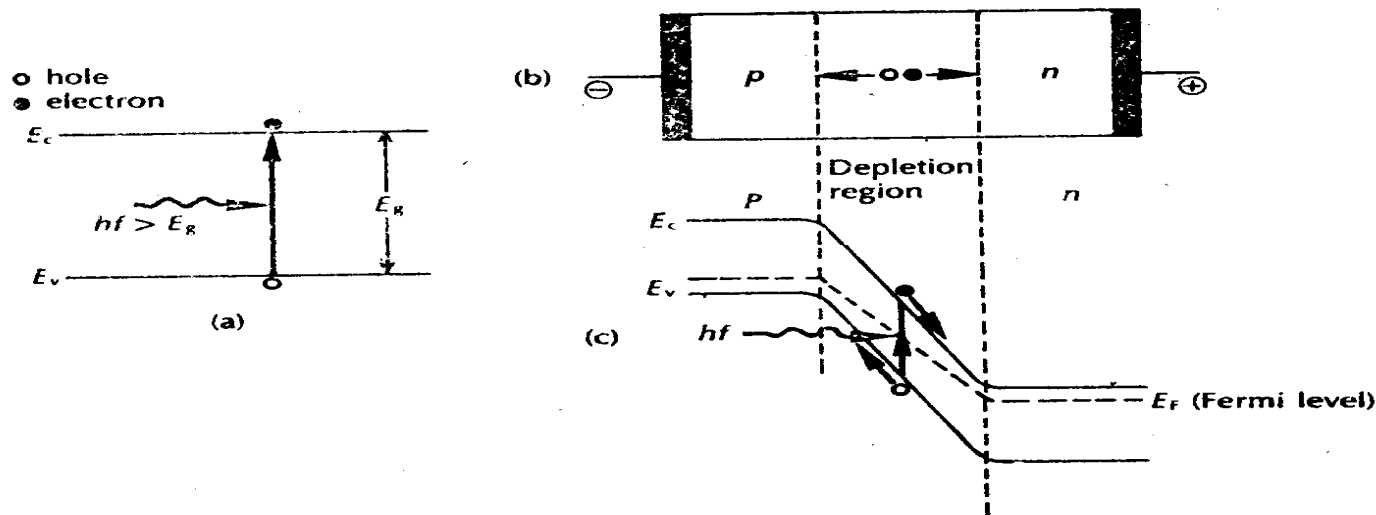


Figure 10.1 Operation of the p - n photodiode: (a) photogeneration of an electron-hole pair in an intrinsic semiconductor; (b) the structure of the reverse biased p - n junction illustrating carrier drift in the depletion region; (c) the energy band diagram of the reverse biased p - n junction showing photo-generation and the subsequent separation of an electron-hole pair.

OPTICAL DETECTION - PRINCIPLE

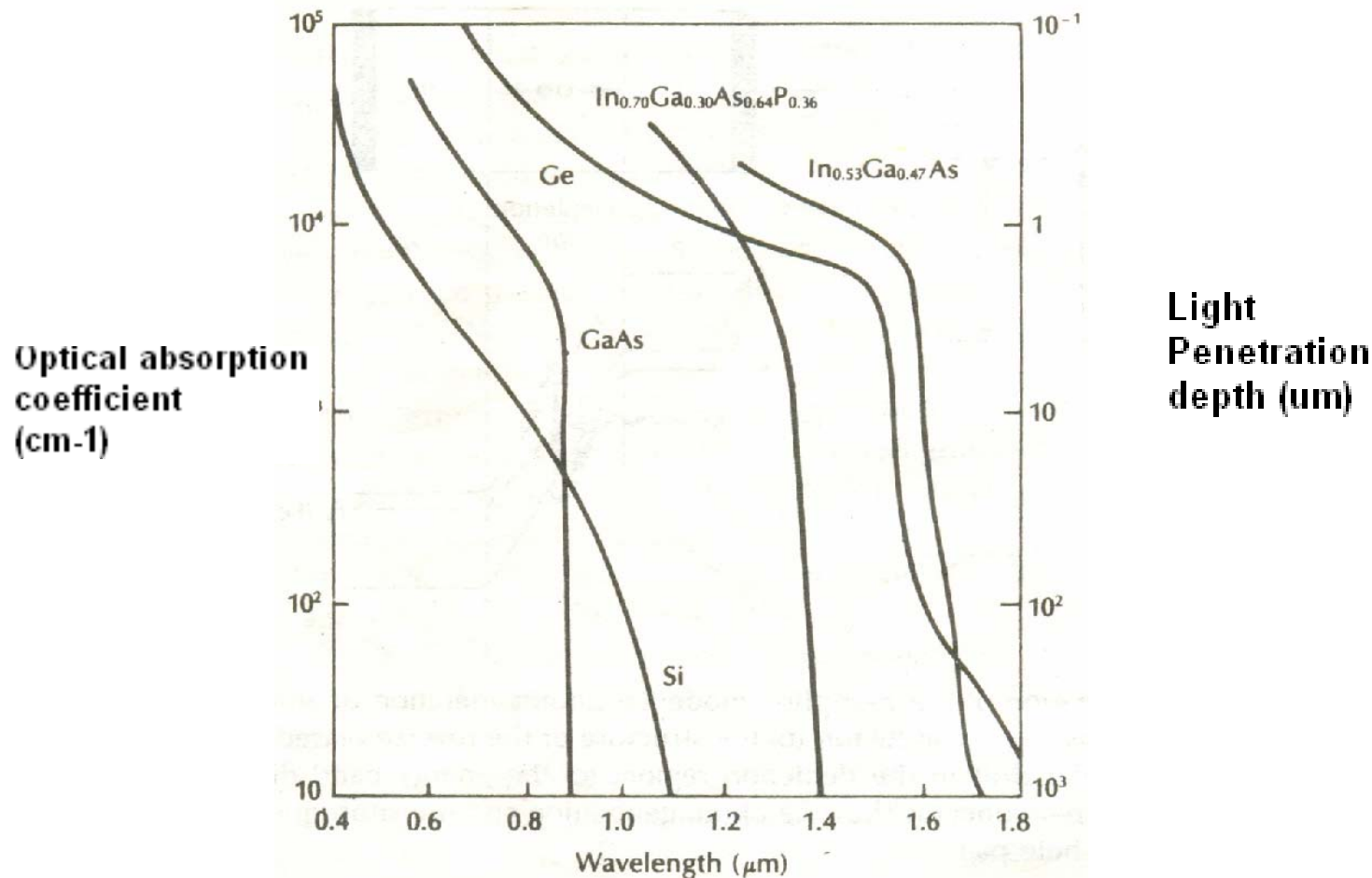
- A Photon incident in or near depletion region which has energy, $hf \geq E_g$ will excite an electron from valence band into conduction band.
- This process creates an electron – hole (carrier) pair.
- Carrier pairs so generated are separated and swept under the influence of electric field . This is known as **displacement current** (in excess of any reverse leakage current)

Opt. Detection-Principle (contd)

- Note: Wider depletion layer (DL) is required for incident light to be absorbed for max. carrier-pair generation.
- However long carrier drift times in DL restrict speed of operation, and hence trade off between sensitivity & response

Absorption coefficient(α_0)

- The absorption of photons in a photo diode to produce carrier pairs and thus a photo current, is dependant on the absorption coefficient, α_0 of the light in the semiconductor (used to fabricate the device)
- I_p (Photo Current) = $P_0 e(1-r) [1 - e^{-\alpha_0 d}] / hf$
- P_0 = Optical power (of incident light)



Optical absorption curves for some common semiconductor: photodiode materials (silicon, germanium, gallium, arsenide, indium gallium arsenide, indium gallium arsenide phosphide)

Band gaps for some semiconductor Photodiode materials at 300k

	Bandgap(ev) at 300k	
	INDIRECT	DIRECT
Si	1.14	4.10
Ge	0.67	0.81
GaAs	-	1.43
InAs	-	0.35
InP	-	1.35
GaSb	-	0.73
InGaAs	-	0.75
GaAs Sb	-	1.15

Absorption coefficient(α_0)

- e- charge of an electron
- r- Fresnel Coefficient (at semiconductor air interface)
- d= width of absorption region.
- Note: **Abs. Coefft. is strongly dependant on λ .** This is due to differing band gap energies of semiconductor photodiode materials at 300 K (Si/Ge/ GaAs/ GaSb/ InAs)

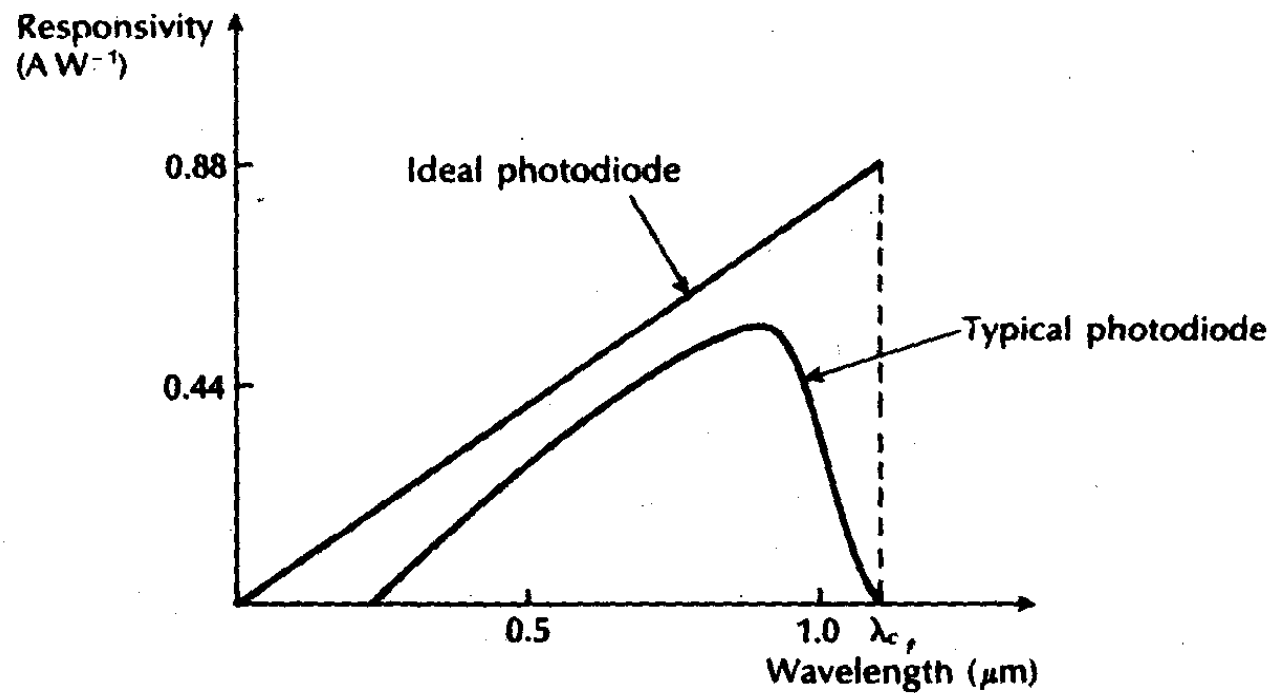
Responsivity (R) of a photo detector

- $R = I_p / P_o$ ($A W^{-1}$)
- I_p = o/p photocurrent (amp)
- P_o = Incident optical power (watts)
- R gives the transfer characteristics of the detector (photo current per unit incident optical power).
- r_p , the incident photon rate $= P_o / hf$ = no.of photons /sec

Responsivity (R) of a photo detector

- But $r_e = \eta r_p$ $[\eta = r_e / r_p] = \frac{\text{No. of electrons collected}}{\text{No. of incident photons}}$
- $r_e = \eta P_o / hf$
- Photocurrent (output), $I_p = \eta P_o e / hf = (r_e * e)$
- $R = I_p / P_o = \eta e / hf$
- $c = f\lambda$ $f = c / \lambda$
- $R = \eta e / hc / \lambda = \eta e \lambda / hc$
- $R \propto \eta$ (at a specific λ)

Responsivity (R) of a photo detector



Long wavelength cut-off

- In intrinsic absorption process ,the energy of incident photons should be greater than or equal to bandgap energy (E_g) of the material of photodetector
- $hf \geq E_g$ or $hc/\lambda \geq E_g$ or $\lambda_c \leq hc/E_g$
- **λ_c =Long λ cutoff (threshold for detection) = hc /E_g**

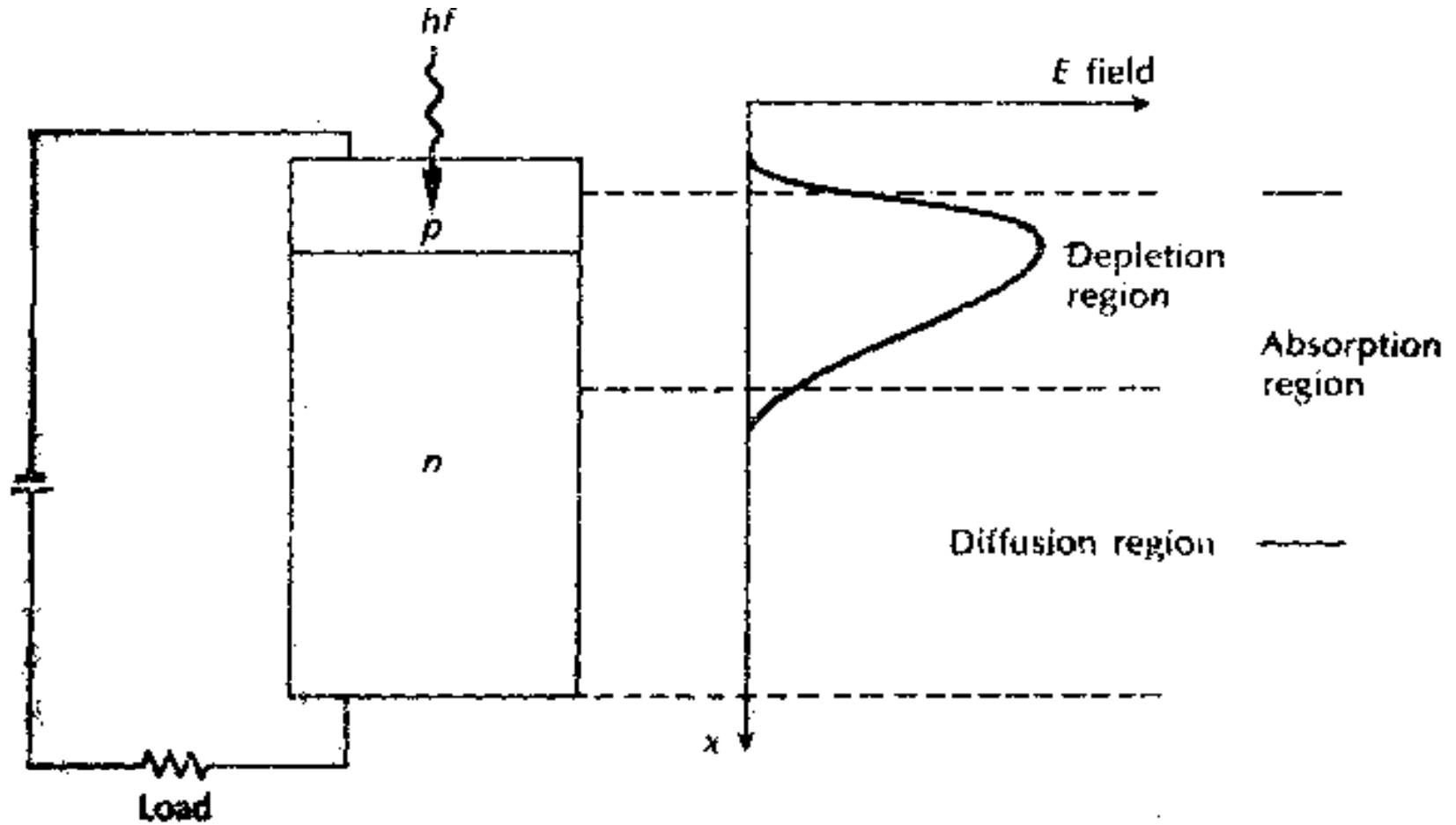
Long wavelength cut-off (contd.)

- This is the longest wavelength of light to give photo detection.
- The cut-off wavelength is $1.06\text{ }\mu\text{m}$ for Si and $1.6\text{ }\mu\text{m}$ for Ge.

Note

- Extrinsic photo detectors are not currently used in OFC.

Reverse bias P-N Photodiode



Reverse bias P-N Photodiode

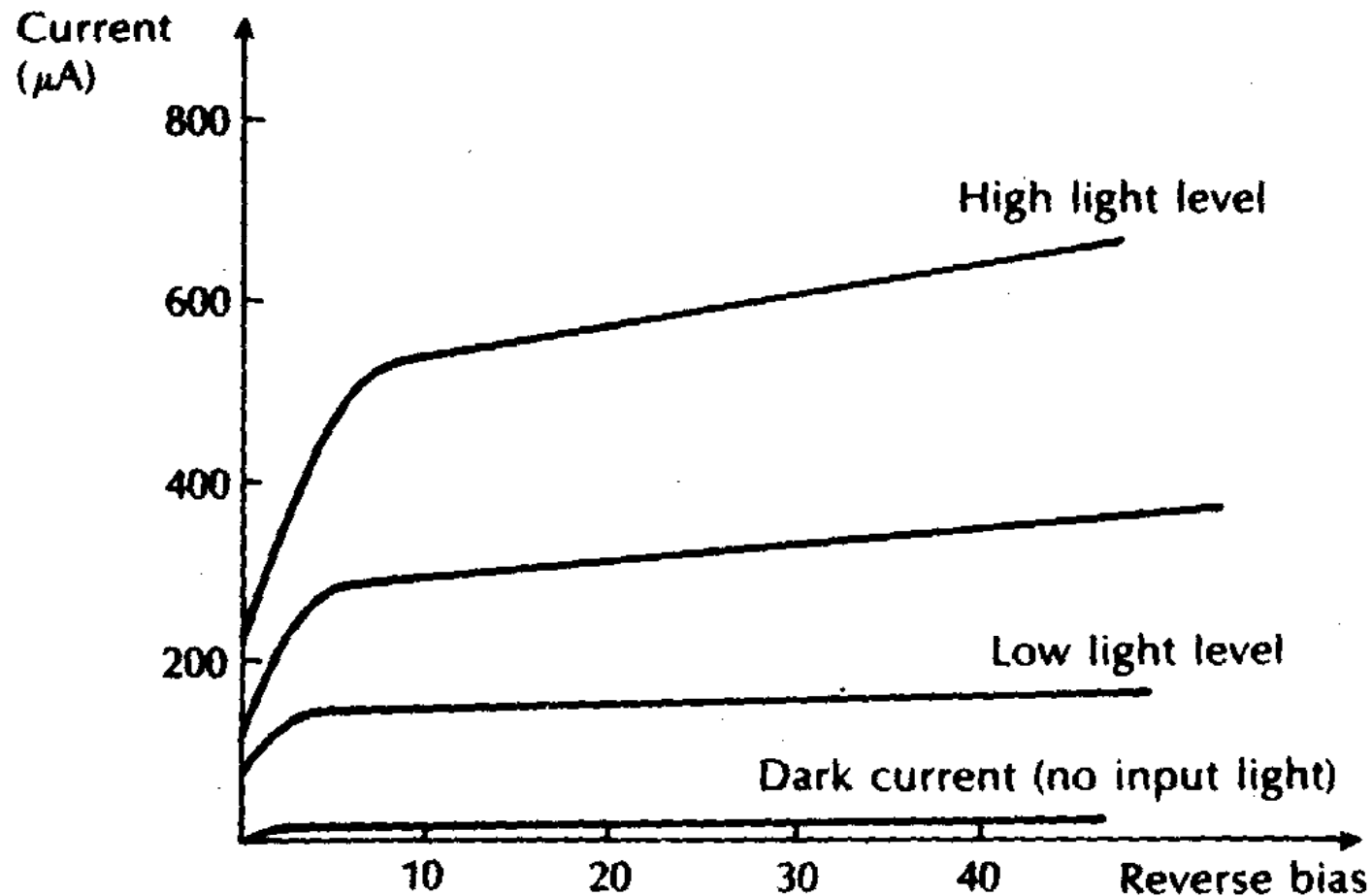
- Width of depletion region is dependant upon doping level.
- Photons are absorbed in depletion as well as diffusion regions.
- Absorption region's dimension depends on energy of incident photons & material of photodiode.
- Electron–Hole pairs are generated in both depletion & diffusion regions.

Reverse bias P-N Photo diode

(contd.)

- **Diffusion is very slow compared to Drift.**
This limits response of photodiode.
- Photons should be absorbed in depletion region (≈ 1 to $3\text{ }\mu\text{m}$)
- Output increases with light level.
- **Note: Lower the doping, wider is the depletion layer.**

Typical P-N Photodiode o/p characteristics

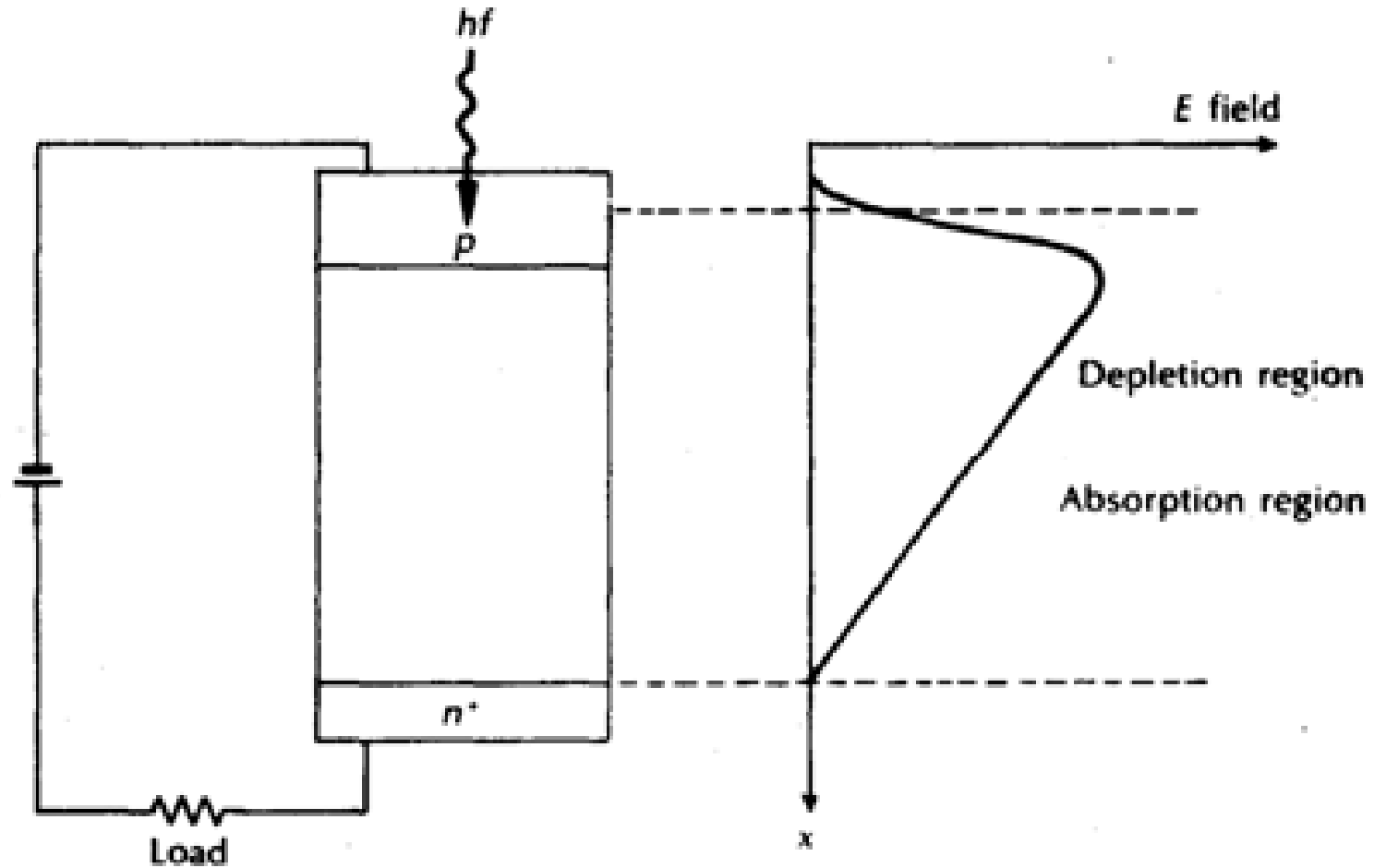


Typical P-N Photodiode o/p characteristics (contd.)

Dark current

- Dark current arises from **surface leakage currents** as well as **generation-recombination currents** in the depletion region in the absence of illumination.
- It may be noted that current o/p is higher for higher light level input to photodiode.

P-I-N Photodiode



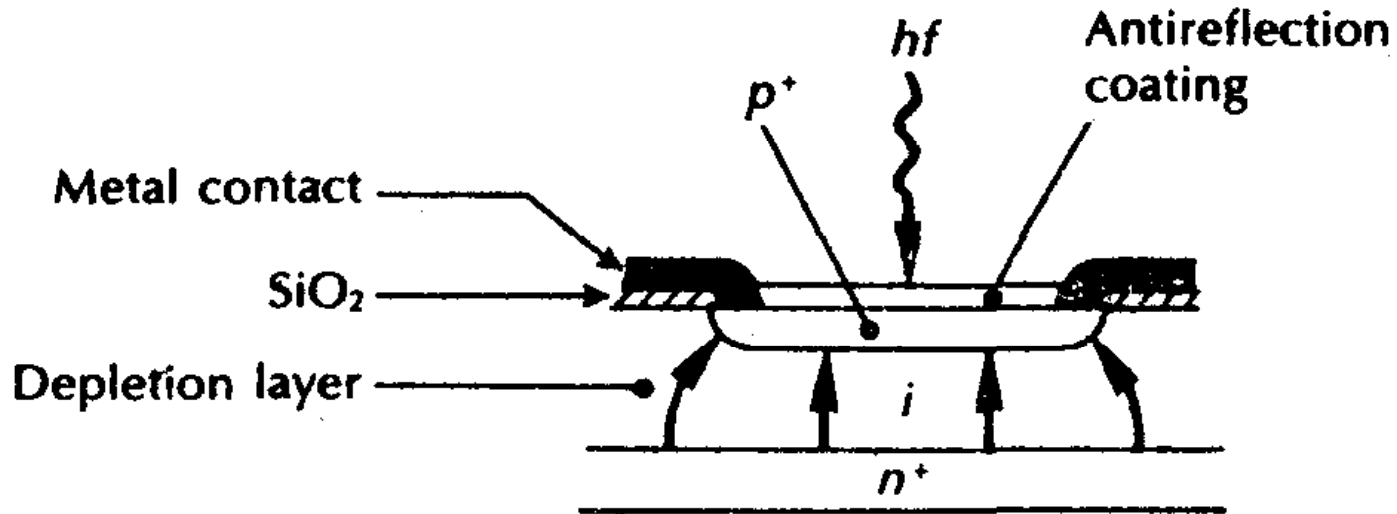
P-I-N Photodiode (contd.)

- **P-I-N STRUCTURE : Highly doped n type (n^+) layer with lightly doped n type (intrinsic) and p type on the other end.**
- A wider depletion layer region is necessary (longer λ)
- All the absorption takes place in the depletion region.

P-I-N Photodiode (contd.)

- Front illuminated photodiode has a fast response time($<1\text{ n sec}$) and low dark current (1 nA). $\eta = 85\%$ $w=20\text{-}50\text{ }\mu\text{m}$ (wavelength= 0.8 to $0.9\text{ }\mu\text{m}$)
- Side illuminated has a larger absorption width ($=500\text{ }\mu\text{m}$) at wavelength= $1.09\text{ }\mu\text{m}$

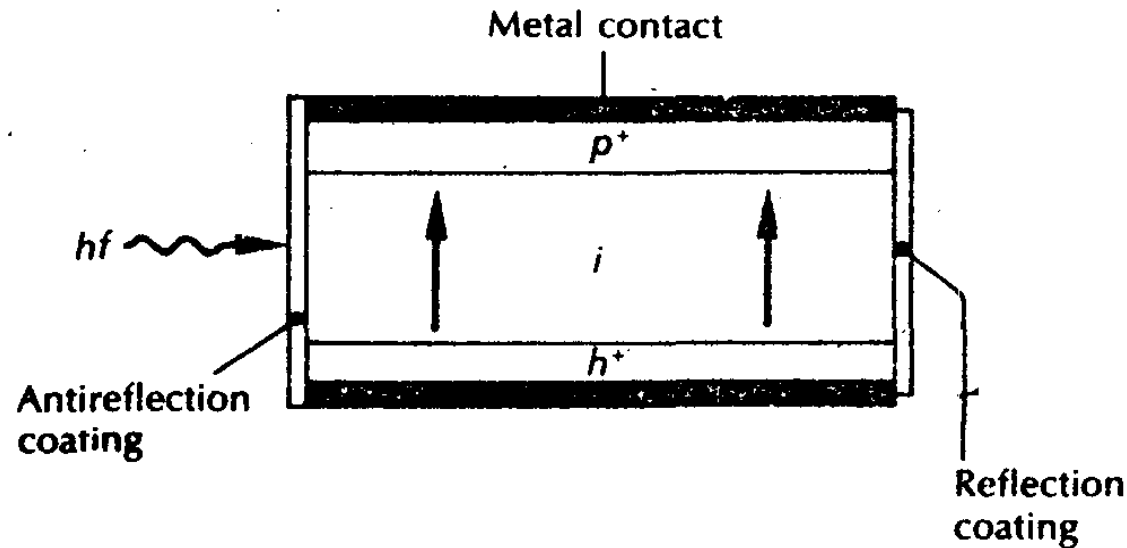
Front illuminated Silicon p-i-n photodiode



Depletion region : 20 to 50 μm for quantum $\eta = 85\%$

- Simplest structure – light entry thr' upper p^+ layer
- **Quantum η is high and dark current is low (1n amp)**
- **Device has fast response time ($< 1\text{ns}$)**

Side illuminated p-i-n photodiode.

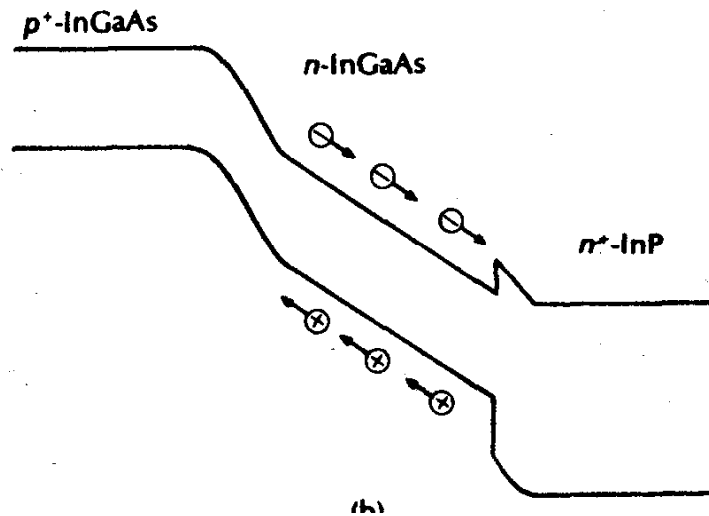
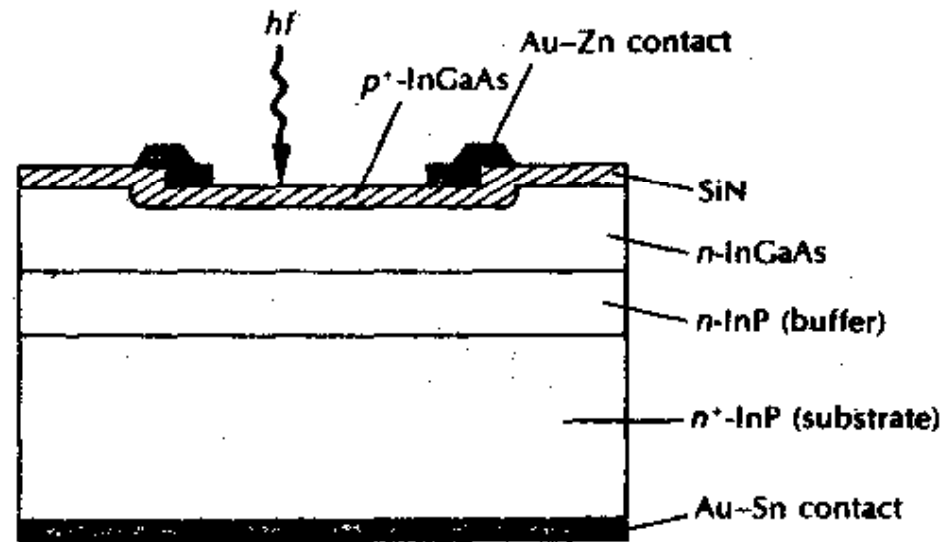


Light is injected parallel to junction plane.

This exhibits large absorption width ($=500\ \mu\text{m}$)

This device is sensitive at wavelength of $1.09\ \mu\text{m}$,
where absorption coeff. is relatively small.

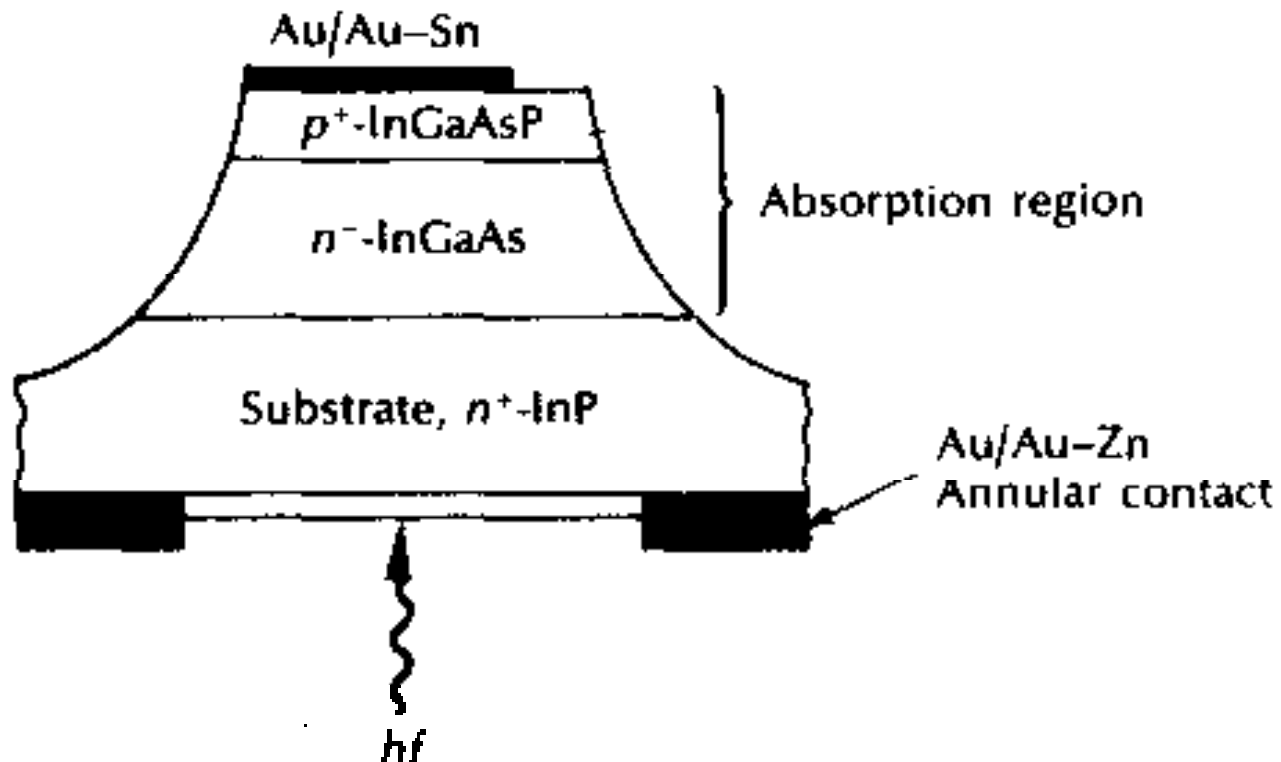
Planar InGaAs p-i-n photodiode-Top entry



SUBSTRATE ENTRY –P-I-N PHOTODIODE

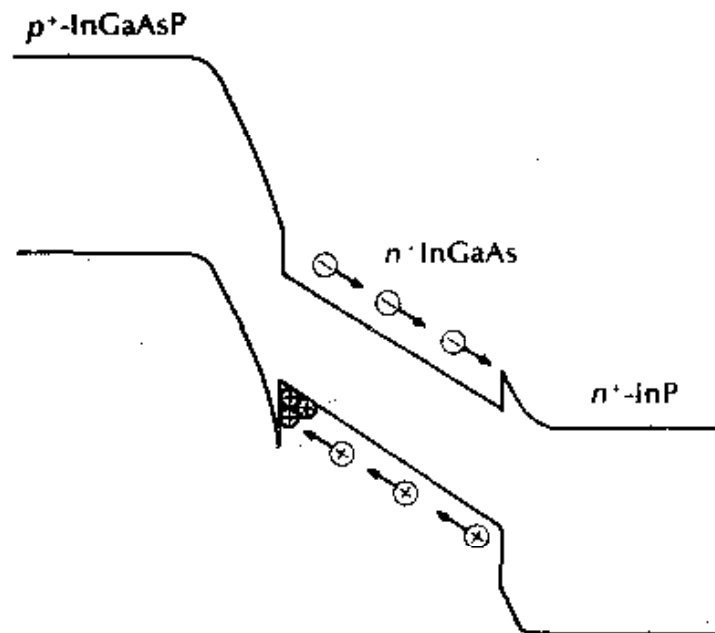
- Fabricated as **mesa structure** which reduces parasitic capacitances.
- Charge trapping at interface causes limitation in response time .

Hetrojunction structure improves quantum η . Such devices can be produced with low capacitance (< 0.1 pF)



Quantum $\eta = 75$ to 100%

Dark current $< 1\text{nA}$



In both devices low doping is used.

BW ≈ 15 GHz (theoretical)

BW = 1 to 2 GHz (typical value practically)

Speed of response (Photodiode)

Factors limiting the speed of response

- a) Drift time of carriers thr' depletion region .
- $t(\text{drift}) = \frac{w}{V_d}$
- Transit time (when $w=10 \mu\text{m}$) = 0.1 n sec.
- Diffusion time of carriers (generated outside w)
- $t_{\text{diff}} = \frac{d^2}{2D_c}$ $d = \text{distance}$
 $D_c = \text{minority carrier diffusion coefft.}$
- $t_{\text{diff}}(\text{holes}) = 40 \text{ nsec.}$ $t_{\text{diff}}(\text{electrons}) = 8 \text{ nsec}$
 $[d=10 \mu\text{m in silicon}]$

Speed of response (contd.)

- Time constant due to cap of photodiode with its load.
- In Cap $C_j = \frac{\epsilon_s A}{\omega}$ ϵ_s = Permittivity of SC.
- ω A-Diode In area
 ω -width of dep. layer
- C_d (cap. of photodiode) = C_j + cap of leads/packaging.
- C_d must be minimised in order to reduce RC time constant which limits the detector response time and bandwidth.

Speed of response (contd.)

- $$B_m \text{ (max 3db BW)} = \frac{1}{2 \pi t_{\text{drift}}} = \frac{1}{2 \pi w/v_d}$$

$$= v_d / 2 \pi w$$

(assuming $c_j=0$, and no carriers are generated outside the depletion region)

- Max. response time of the device $= 1/B_m$

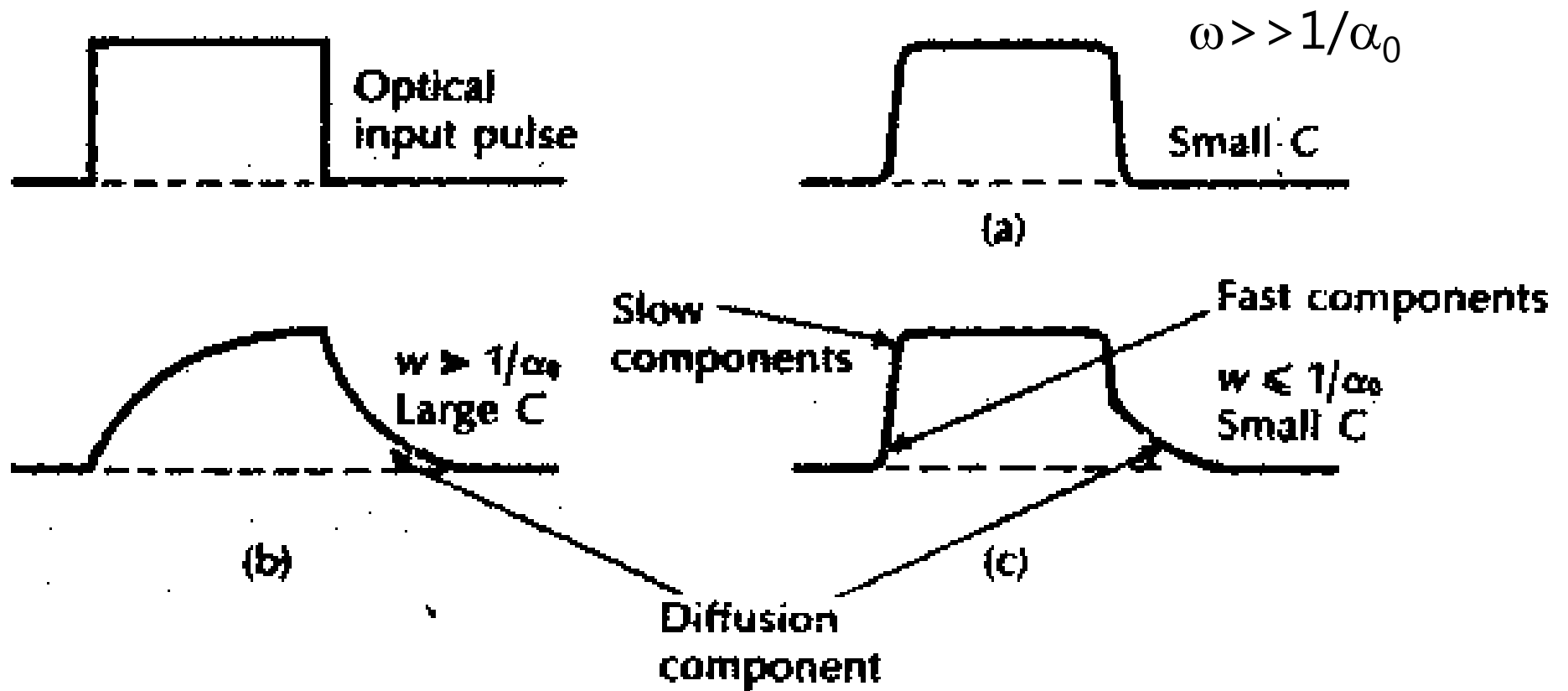
Response of photodiode (Contd.)

- **For higher quantum η , $w \gg 1/\alpha_0$ so that most of the incident light will be absorbed.**
- **Fig A : Response of photodiode with the above condition and small capacitance (negligible diffusion outside depletion region).**
- **Fig B: (Large C) Speed of response becomes limited by the RC time constant (R-Load Res.)**

Response of photodiode (Contd)

- Fig C: ($w < 1/\alpha_0$) and (small C) **Dep. layer is narrow.** Carriers are created by absorption outside the dep. region. o/p pulse displays a long tail caused by the diffusion component.
- Devices with **thin dep. layer** have a tendency to exhibit distinctive **fast response** and **slow response components**. The fast response is from absorption in the thin dep. layer

Response of photodiode (Contd.)



Noise in photodiodes

- Random current and voltage fluctuations occur at o/p in the presence as well as absence of an optical signal.
- **Dark Current (I_d) : O/P photocurrent in the absence of an opt. input signal.**
- RMS value of shot noise *current*

(A)

$$(\overline{i_s^2})^{1/2} = (2e\overline{BI})^{1/2}$$

Noise in photodiodes (Contd.)

- where $\overline{i_s^2}$ = mean square current variation
 B = Rx Bandwidth
 I = Detector average current.
- Noise performance is assessed using the following
 - **Noise Eqvt. Power (NEP)**
 - **Detectivity (D)**
 - **Specific Detectivity (D*)**

Noise in photodiodes (Contd.)

- **NEP: It is the incident opt power at a particular λ required to produce a photo current equal to RMS noise current within a unit bandwidth**

From Responsivity, $I_p = \eta P_0 e / hf$ or $P_0 = I_p hf / \eta e$

Or $P_0 = I_p h c / \eta e \lambda$,

using \overline{A} and putting $I_p = (\overline{i_s^2})^{1/2}$

$$I_p = (2e B I)^{1/2} \quad \text{where } I = I_p + I_d, \quad I_p = [2e B (I_p + I_d)]^{1/2}$$

$$\text{when } (I_p \gg I_d), \quad I_p = 2eB$$

$$\text{NEP} = P_0 = 2eB \cdot h c / (\eta e \lambda) = 2hc / \eta \lambda$$

Noise in photodiodes (contd.)

- When $I_p \ll I_d$ (dark current)
- Then $I_p = [2e(I_p + I_d)B]^{1/2} = [2eI_d B]^{1/2}$
- $I_p = [2eI_d]^{1/2}$ When $B = 1$ Hz
- $\text{NEP} = P_0 = \frac{hcI_p}{\eta e\lambda} = \frac{hc(2eI_d)^{1/2}}{\eta e\lambda}$
- **Detectivity (D)** = $1/\text{NEP} = \frac{\eta e\lambda}{hc(2eI_d)^{1/2}}$

Noise in photodiodes (Contd.)

- **Specific Detectivity (D^*)** :The area (A) of Photodetector is taken into account. This is necessary when background radiation and thermal generation rather than surface conduction are the major causes of dark current .

$$\begin{aligned}
 \bullet \quad D^* &= DA^{1/2} = \frac{\eta e \lambda}{hc(2eI_d)^{1/2}} \quad (A^{1/2}) \\
 &= \frac{\eta e \lambda}{hc(2eI_d/A)^{1/2}} \quad [\text{when } B=1 \text{ Hz}]
 \end{aligned}$$

$$D^*(\text{over } BW=B) = D(AB)^{1/2}$$

Noise (contd.)

- **Thermal Noise** :This is the spontaneous fluctuation due to **thermal interaction between the free electrons and vibrating ions** in a conducting medium (especially prevalent in resistors at room temp.)
- Thermal noise current $\overline{i_t^2} = \frac{4 KTB}{R}$
- where K is Boltzmann's Constant
B-Post Detection (Elect) BW
R- Resistor in optical Receiver

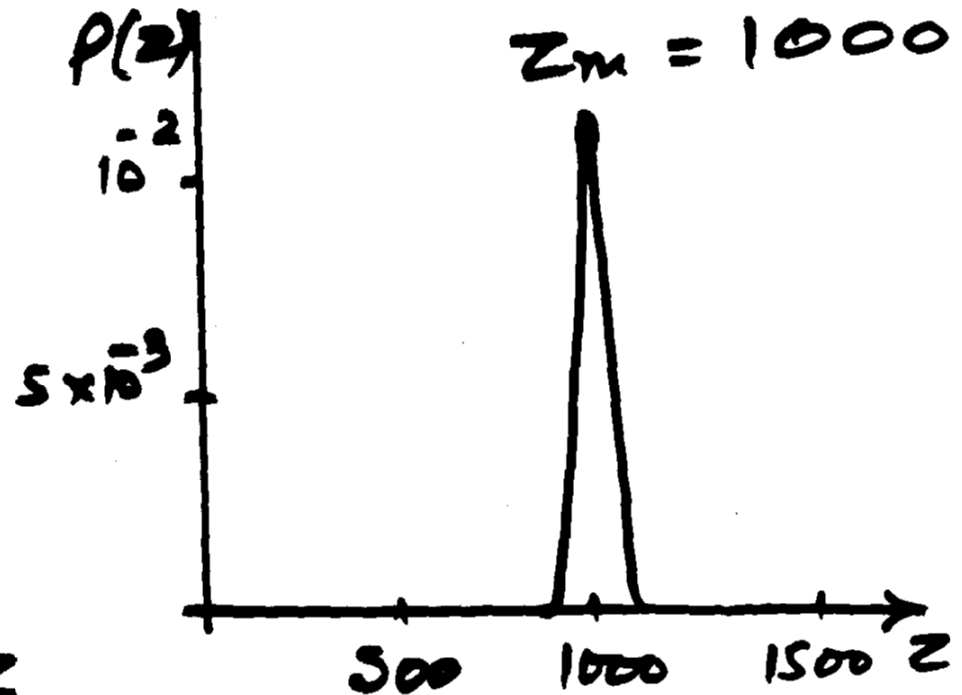
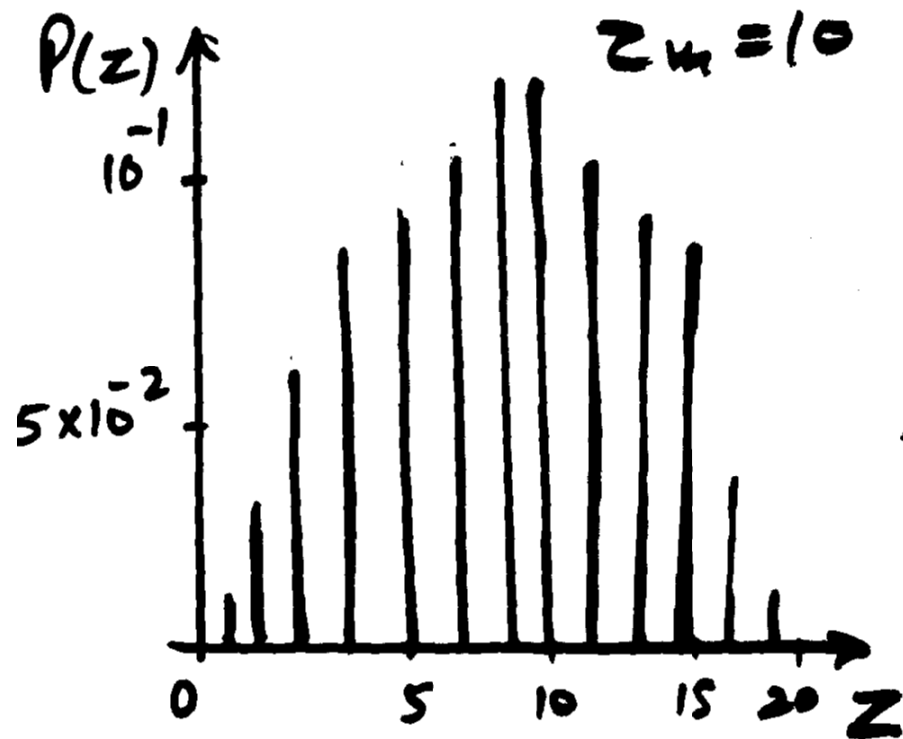
Noise (contd.)

- **Dark Current Noise**
- $i_{(d)}^2 = 2eBI_d$
- Dark current can be reduced by careful design and fabrication of the detector.
- **Quantum Noise**
- $E=hf$ (energy of a photon)
- At optical frequencies, $hf > KT$, therefore the quantum behaviour of e.m. radiation must be taken into account as quantum fluctuations dominate over thermal fluctuations.

Noise (Contd.)

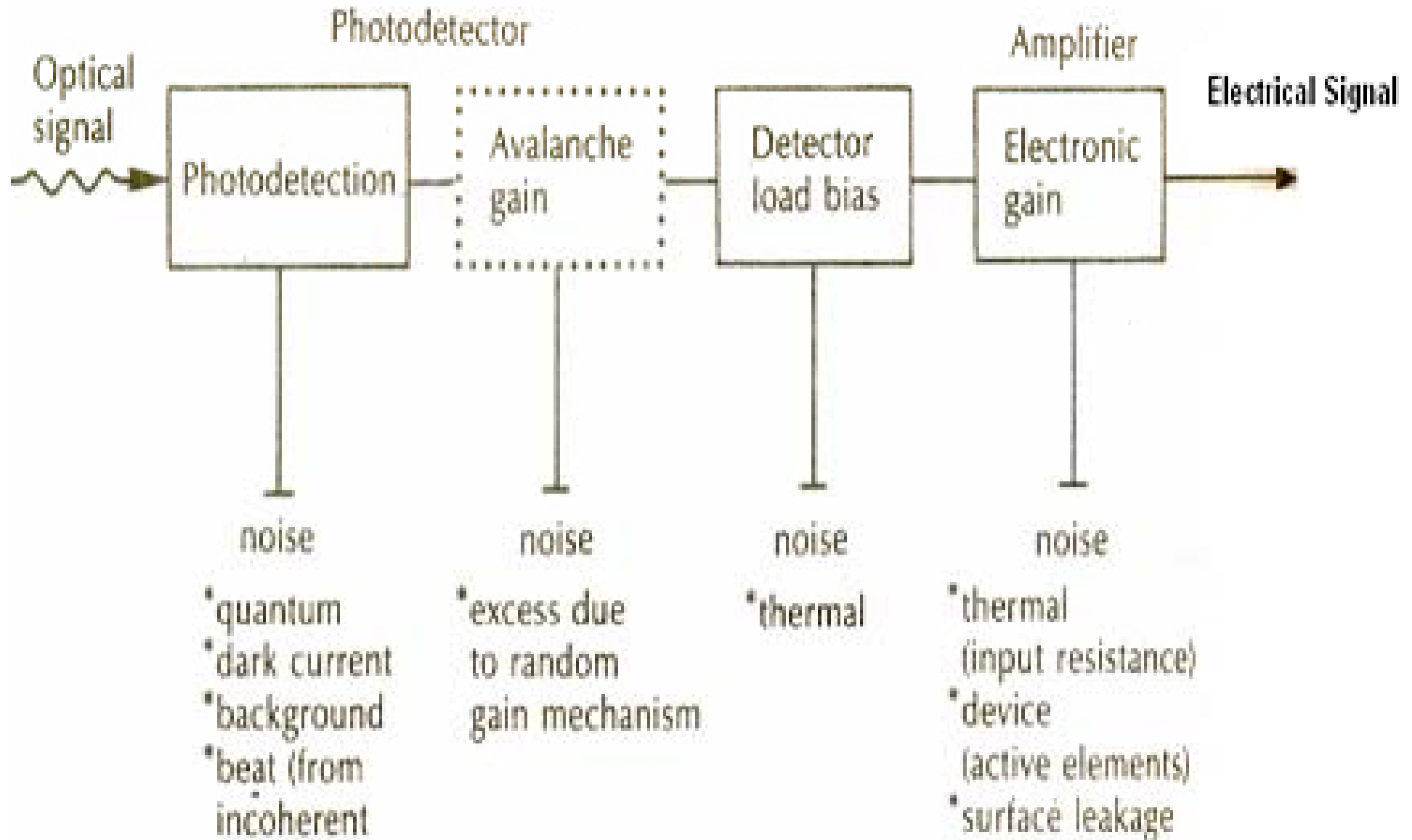
- $P(z) = \frac{z_m^z e^{-z_m}}{z!}$ Where z_m = Variance of prop. distribution
- $P(z)$ = Prob of detecting z photons in time period t
- Z_m = mean = Variance (Poisson Distribution)
- $\eta = \frac{r_e \text{ (electron rate)}}{r_p \text{ (incident photon rate)}}$
- or $r_e = \eta r_p = \eta \frac{P_o}{hf}$
- No. of electrons generated in time t is equal to avg. no of photons detected over this time period.
- $Z_m = \frac{\eta P_o}{hf} t$

Poisson Dist for $Z_m = 10$ & $Z_m = 1000$



This represents the detection process for monochromatic coherent light.

NOISE IN OPT. RECEIVER



APD Design

– Desirable Features

- Carrier multiplication should take place uniformly across the whole area illuminated by the incident radiation.
- The peak field where avalanche multiplication will occur should be confined to a very thin layer.
- The avalanche should be initiated by carriers with higher ionisation coefficient, otherwise BW will be less & noise factor will be increased.
- High quality material should be used (to prevent premature avalanche)

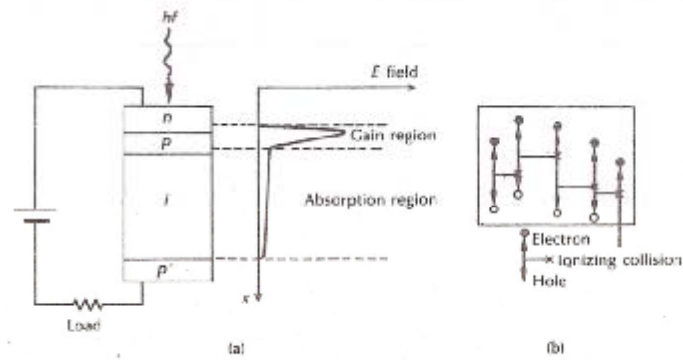
APD Design (Contd.)

- Quantum η should be high
- Dark currents should be low.
- (Note: Higher dark currents are due to small energy band gap and edge/surface defects)
- Noise Factor should be low
- Speed of response should be high.

APD Design (contd.)

- Note: (The speed of response of a photodiode is limited by the time it takes the photo-generated carriers to drift across the depletion region)
- The system design should take into account various impairments (modal noise, dispersion, feedback noise, cross talk and non-linearities in fiber.)

APD



(a) Avalanche photodiode showing high electric field (gain) region. (b) Carrier pair multiplication in the gain region.

APD

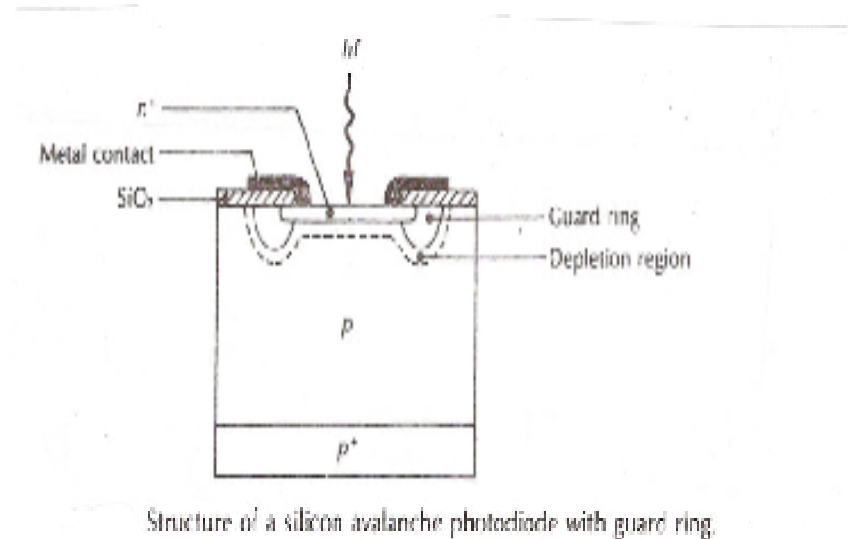
- More sophisticated structure than p-i-n Photodiode so as to create a high electric field.
- Most of the photons are absorbed in the Depletion region .
- There is another high field region in which holes and electrons acquire sufficient energy to excite new electron-hole pairs.(IMPACT IONISATION)

APD (contd)

- High Reverse Bias required (50-400V)
- Carrier multiplication factors upto 10^4 can be obtained using defect free materials
- **MICROPLASMAS**(small areas with lower BD voltage than the remainder of the junction) must be reduced thr better fabrication and defect free materials.
- Rise time: 150-200 ps, Fall time: 1ns
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SILICON APD(with guard ring)

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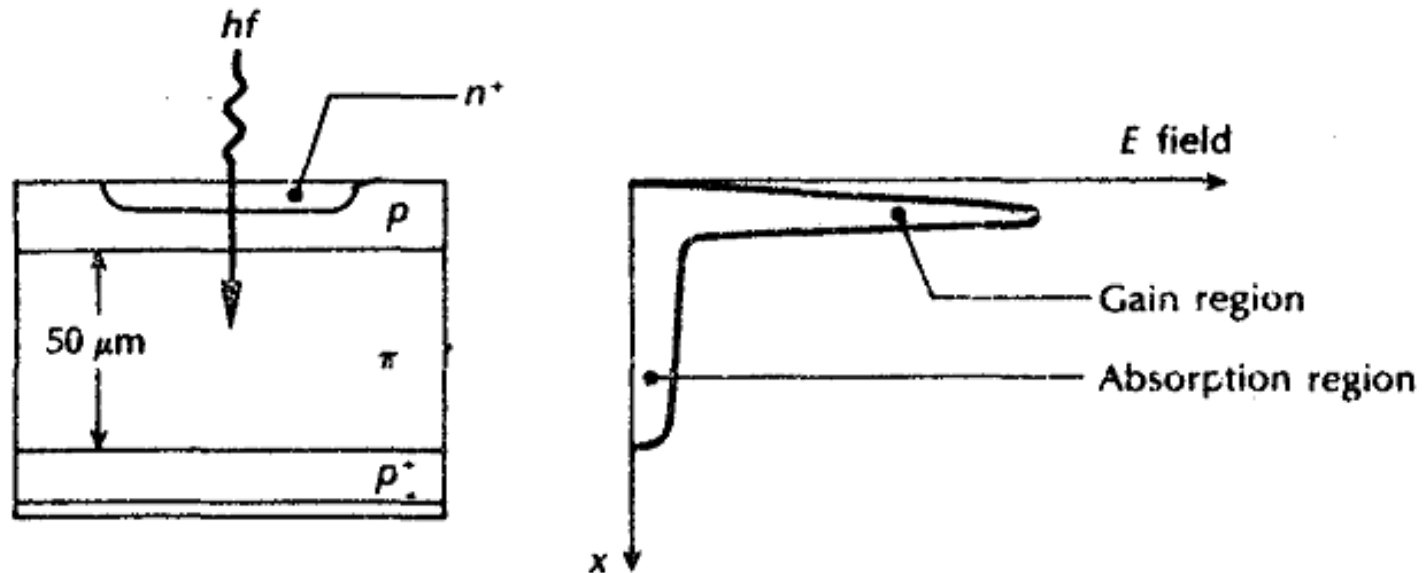


SILICON APD

- Guard Ring structure eliminates excessive leakage at junction edges.
- When elect field is high, all carriers drift .
- At low gain, transit time and RC time constant dominate .
- At high gain, Avalanche build up time dominates (BW decreases with increasing gain)

Silicon Reach Through APD(RAPD)

- For min noise, the elect field at avalanche breakdown must be as low as possible and impact ionisation should be initiated by electrons.
- RAPD consists of $p^+ - \Pi - p - n^+$ layers
- Avalanche multiplication takes place in relatively narrow place centered on $p - n^+$ junction.



Silicon Reach Through APD(RAPD)

- When Reverse Bias increases, depletion layer widens until it reaches through to lightly doped Π region

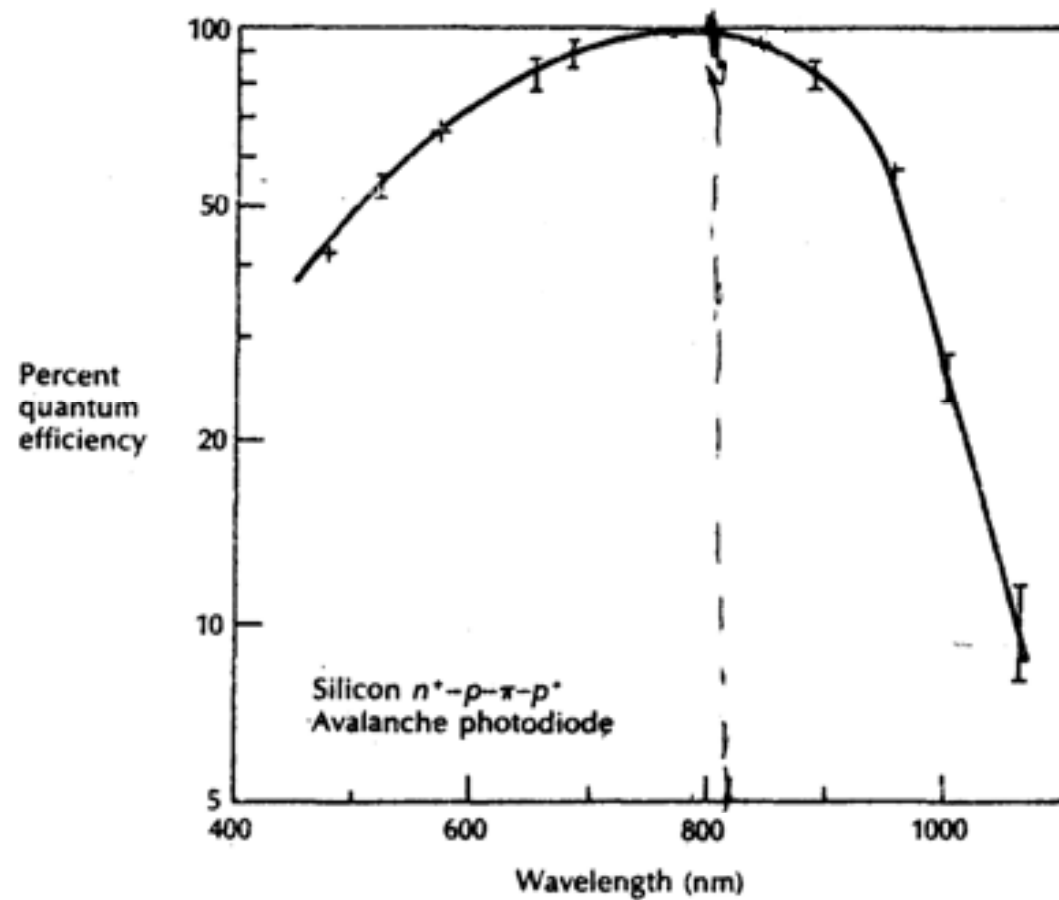
Field in Π region is much lower than at $p-n^+$ junction, but still it is high enough (2×10^4 V/cm) .This limits the transit time and ensures a fast response (0.5 n sec.)

- At $0.825 \mu\text{m}$ wavelength, Quant $\eta = 100\%$
- Dark currents are low.

RAPD (contd)

- Normally RAPD is operated in fully depleted mode
- Light enters the device thr p^+ region and is absorbed in the π material
- Upon being absorbed, photon gives up its energy, thereby creating electron-hole pairs
- The electrons drift thr π region in the pn^+ junction, where high electric field exists
- Carrier multiplication takes place in this high-field region.

Quantum Efficiency Vs wavelength Silicon RAPD



Quantum Efficiency Vs Wavelength Silicon RAPD

- Quantum = 100% at 0.825 μm (λ)
- Dark Currents for this photodiode are low and depend only slightly on bias voltage.

Multiplication Factor(M) –APD

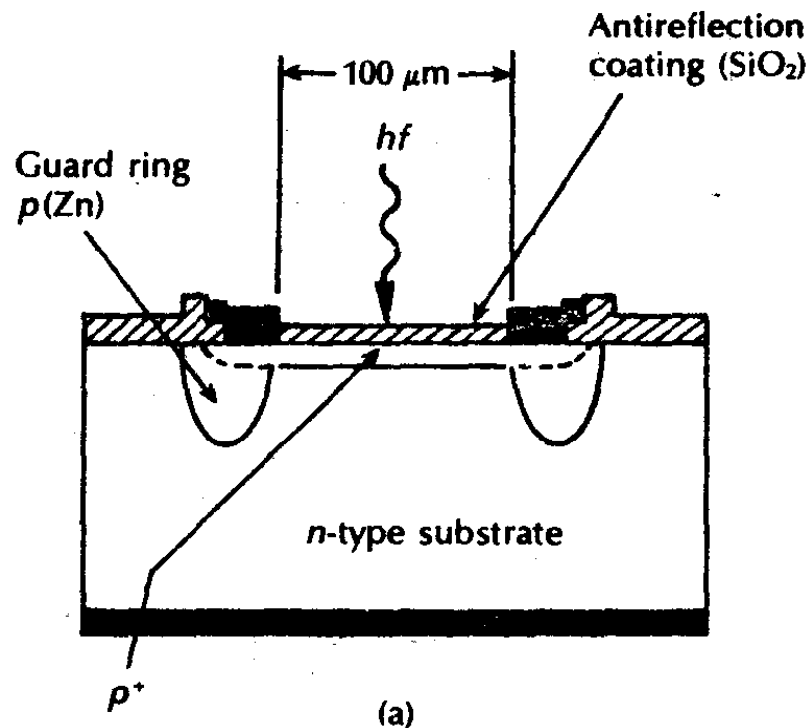
- M is a measure of the internal gain provided by the APD.
- $M = I / I_p$
- Where I - Total O/P current at the operating voltage (where carrier multiplication occurs)
- I_p - Initial or primary photo current (i.e. before carrier multiplication occurs)

Germanium Avalanche Photodiode

- These APD's can be used over the entire wavelength range (0.8-1.6 μm).
- However ,these APD's have **higher dark currents** together with **excess noise** than those in silicon APD's
- These higher dark currents are due to surface effects and excess noise is due to electrons rather than holes initiating the multiplication process.

Germanium Avalanche Photodiode (contd)

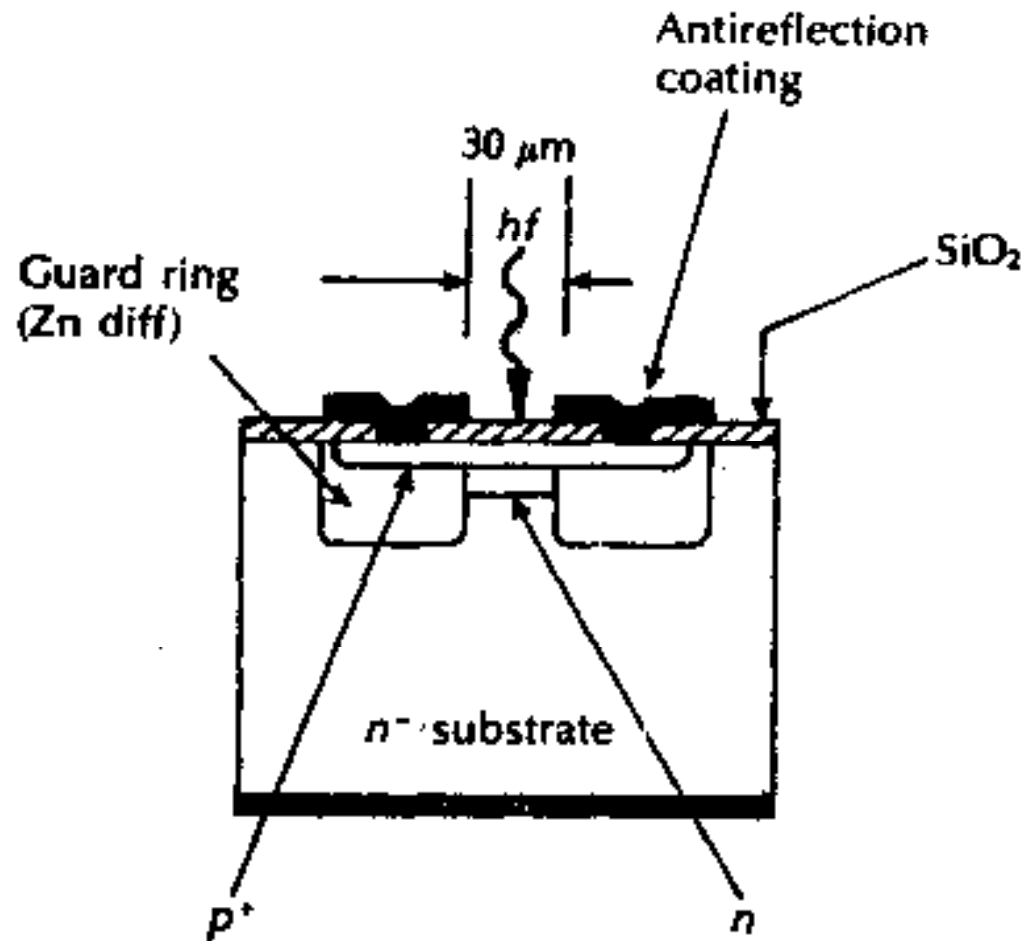
- Advantage: Avalanche BD voltage is low (typically 25 v)
- Ge has a higher absorption co-efficient



- **Ge APD – $p^+ n$ structure**
- **Dark current**

150 to 250 nA	100 μm sensitive area
5 nA	30 μm sensitive area
- Speed of this structure at 1.5 μm wavelength is poor.(most of abs. outside DL)
- This is overcome by the use of $p^+ n n^-$ structure, also known as Hi –LO structure.

Ge APD-Hi-LO (p^+nn^-) structure.



Ge APD-Hi –LO ($p^+ nn^-$) structure.

- The above structure resembles the reach thr structure for silicon APD's
- It is called Hi-Lo structure as it combines high bandwidth (700MHz) with low multiplied dark current (33 nA)
- It has good excess noise performance.

Ge APD- Hi-LO (p^+nn^-)structure (contd.)

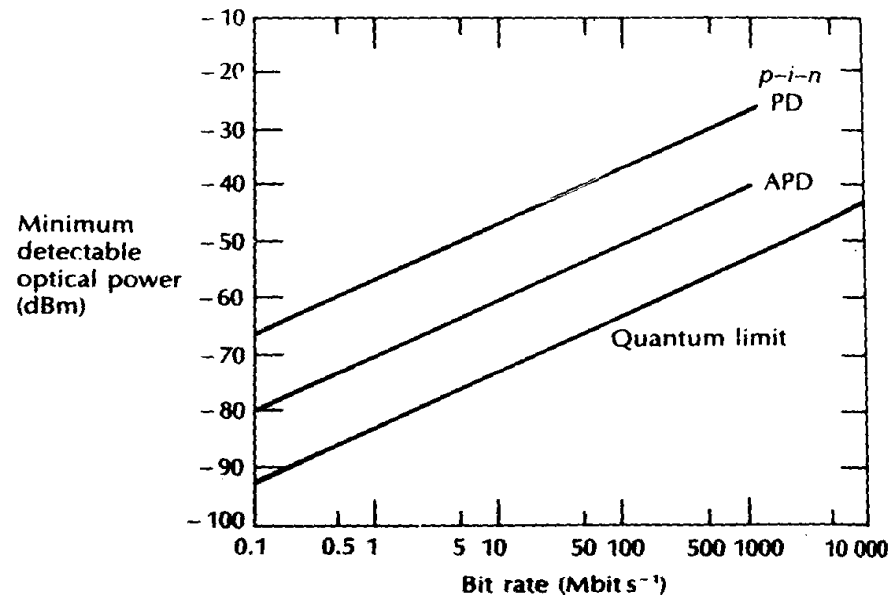
- **Disadvantage:** Breakdown voltage is higher (85V)
- The Hi-Lo devices are among the highest performing Ge APD's for longer wavelength operation.

Benefits & Drawbacks-APD

Si Det.

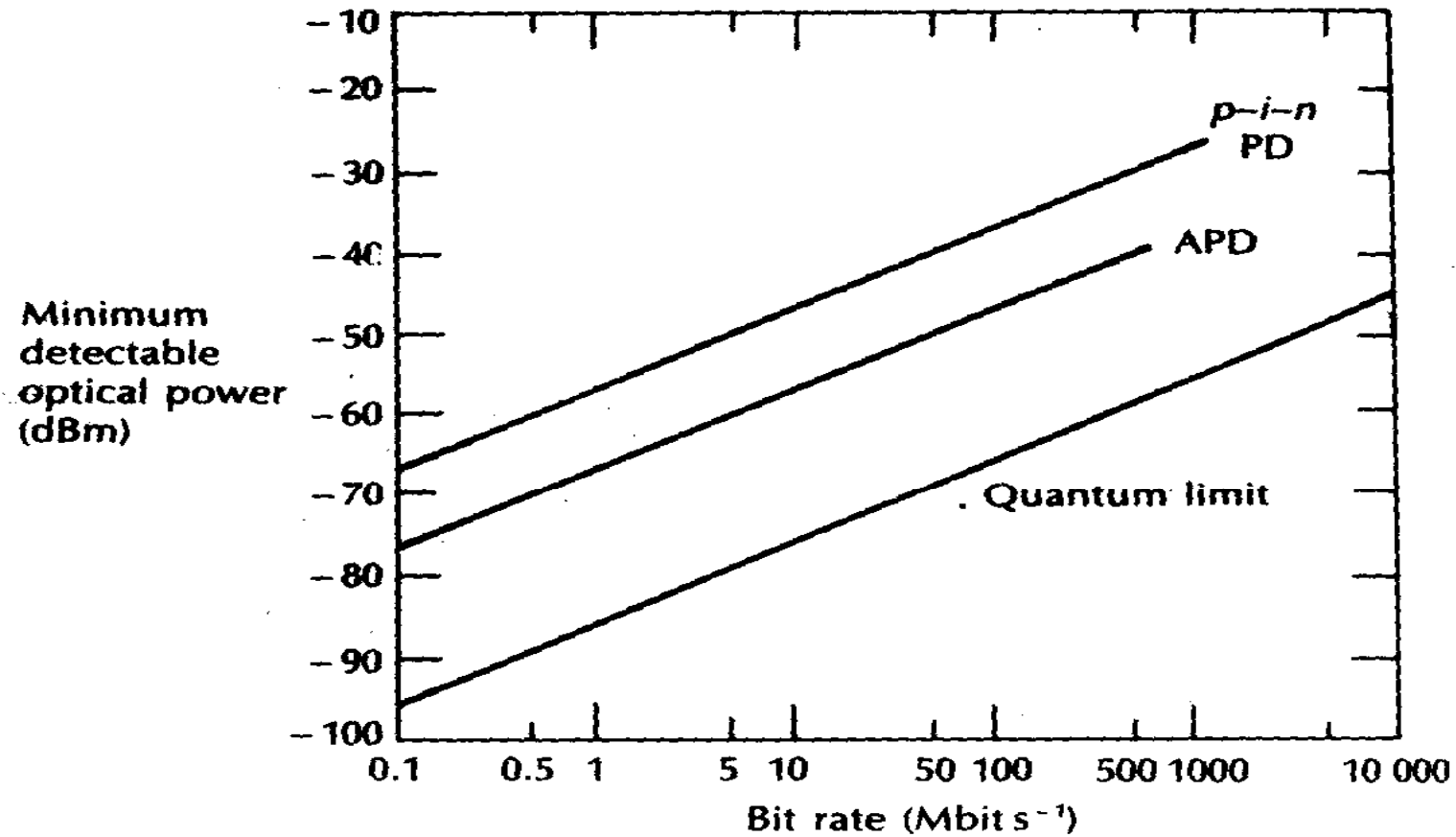
sensitivity at

0.82 μm



Benefits : APD's can detect low light levels

- Increased sensitivity over p-i-n diodes (5 to 15db)
- Wider Dynamic Range (as a result of gain variation with response time & rev. bias)



In GaAs DET at $\lambda=1.55\mu\text{m}$

Drawbacks of APD's

- Fabrication difficulty (complex structure)
- Increased cost.
- High bias voltages (50 to 400V)
- Additional Noise (due to random nature of gain mechanism)
- Variation of gain with temp. This requires a temp compensation for stability of operation.